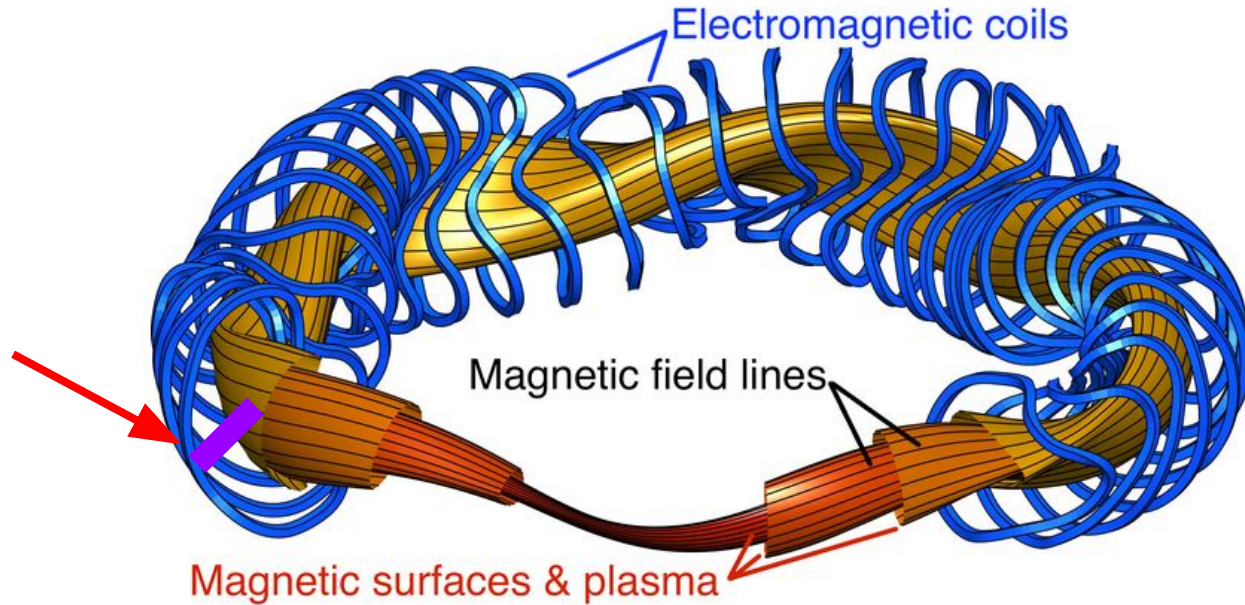


Magnetic Gradient Scale Length

An Analysis of Why Certain Plasmas Require
Close External Magnetic Coils



Pre-print:



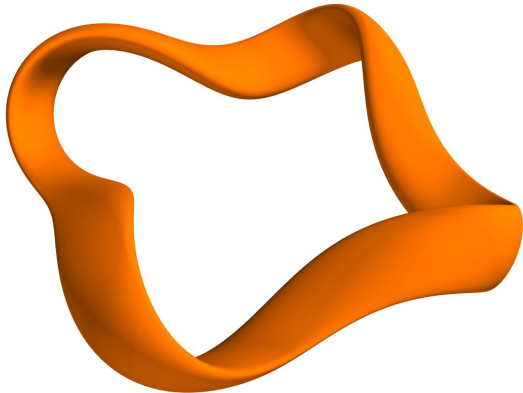
By John Kappel

PI: Dr. Matt Landreman

Optimization of Stellarators is a Two Stage Process

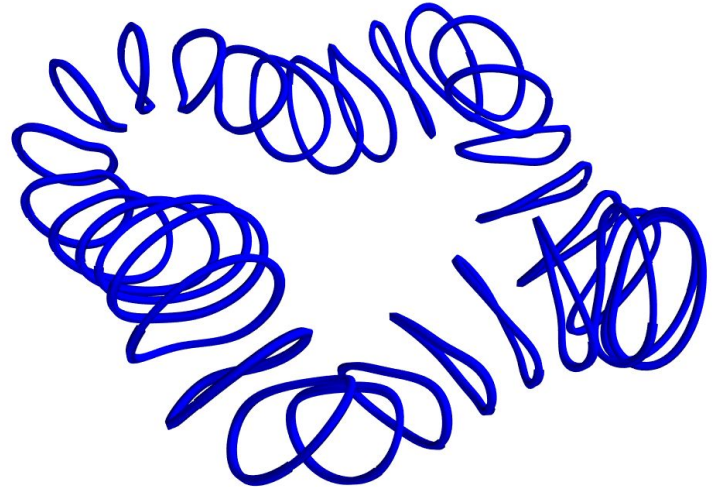
Stage I Optimization

- Last Closed Flux Surface (LCFS) Shape
- Pressure Profile
- Rotational Transform



Stage II Optimization

- Coil Shape
- Subject to Engineering Constraints

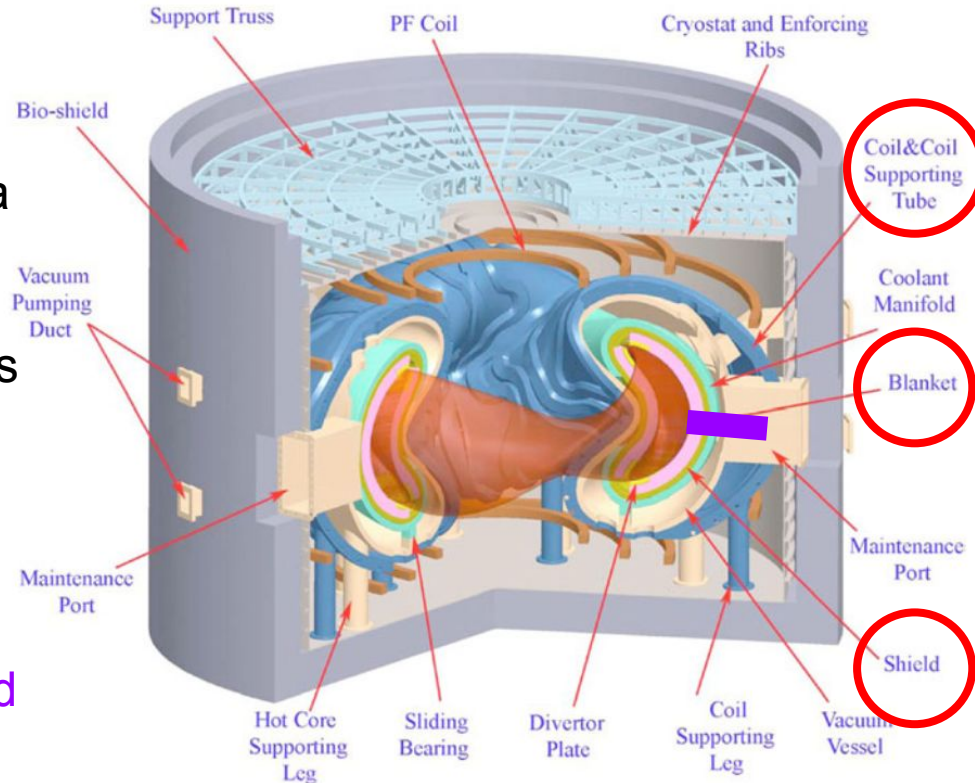


Stellarators Need Space for Breeding Blankets & Neutron Shielding

Need ~1.5 m between coils and plasma

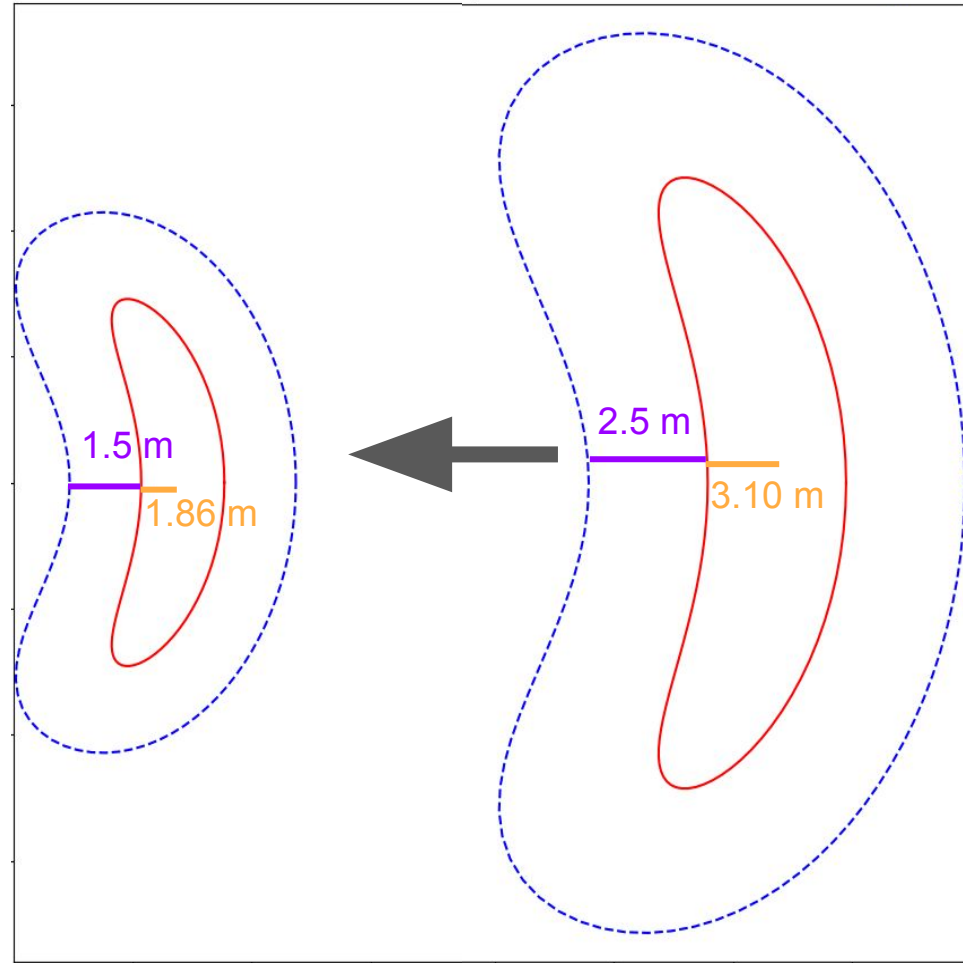
Both ARIES-CS and W7-X report issues with not enough spacing

Plasma-Coil Separation: distance between the last closed flux surface and the center of an external field coil



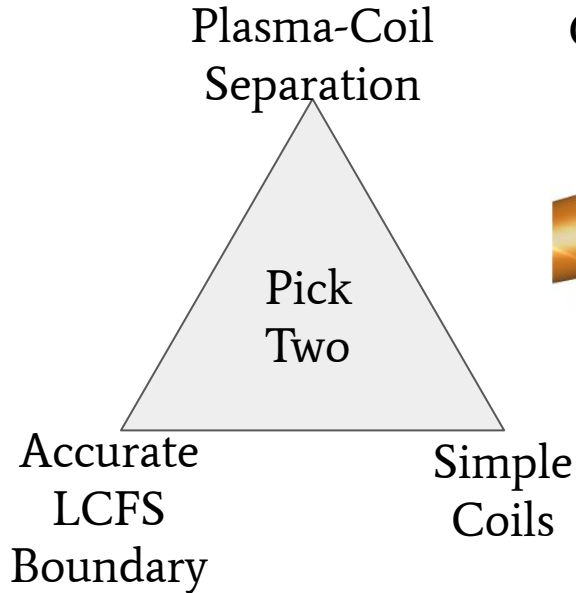
Other Benefits of Large Plasma-Coil Separation

1. Reduced coil ripple
2. Components can shift during startup and initialization
3. Plasma configurations large
plasma-coil separation can be scaled down

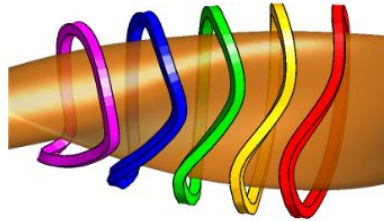


* Precise QH, to scale

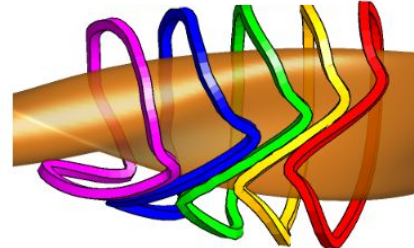
Difficulty of Increasing Plasma-Coil Separation in Stage II Optimization



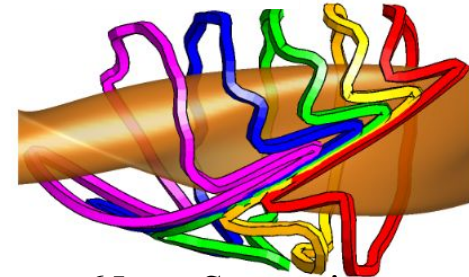
Coils offset a uniform distance from the W7-X plasma



25 cm Separation



50 cm Separation

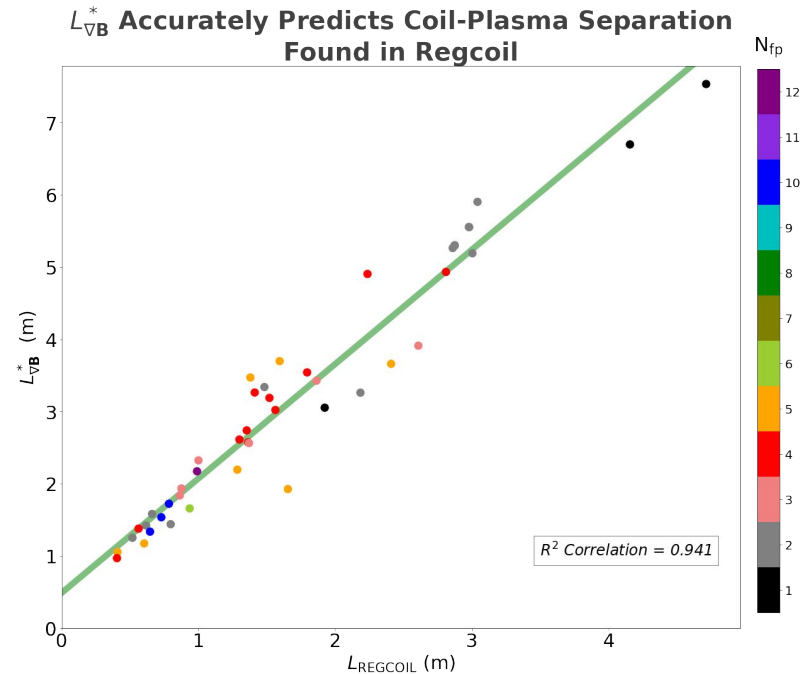


65 cm Separation

Coils that have a higher uniform separation from the plasma are more complex

Hypothesis: Plasma-Coil Separation can be Understood in Terms of Magnetic Gradient Scale Length

We shall show a **good correlation** between this magnetic gradient scale length ($L_{\nabla B}$) and the plasma-coil separations of actual coil designs of over 40 configurations calculated in REGCOIL.



Outline

1. Intuition for Magnetic Gradient Scale Length
2. Methods of Coil Optimization in REGCOIL
3. Comparison Between $L_{\nabla\mathbf{B}}$ and L_{REGCOIL} and Discussion

Outline

1. **Intuition for Magnetic Gradient Scale Length**
2. Methods of Coil Optimization in `REGCOIL`
3. Comparison Between $L_{\nabla B}$ and L_{REGCOIL} and Discussion

Precedent for the Magnetic Gradient Scale Length

The diagram illustrates the derivation of the magnetic gradient scale length. It features three mathematical expressions arranged in a flow from left to right. On the left, there are two scale lengths: L_{Te} and L_n . An arrow points from L_{Te} to the expression $\frac{T_e}{\|\nabla T_e\|}$. Another arrow points from L_n to the expression $\frac{n}{\|\nabla n\|}$. A third arrow points from the $\frac{T_e}{\|\nabla T_e\|}$ expression to the final expression $\frac{Q}{\|\nabla Q\|}$. A fourth arrow points from the $\frac{n}{\|\nabla n\|}$ expression to the same final expression $\frac{Q}{\|\nabla Q\|}$.

$$L_{Te} \rightarrow \frac{T_e}{\|\nabla T_e\|}$$
$$L_n \rightarrow \frac{n}{\|\nabla n\|}$$
$$\frac{Q}{\|\nabla Q\|}$$

Arguments of scale lengths are used in plasma physics to determine which effects are negligible and significant.

A spatial gradient of the magnetic field encodes some information about the spatial distance from the coils to the plasma.

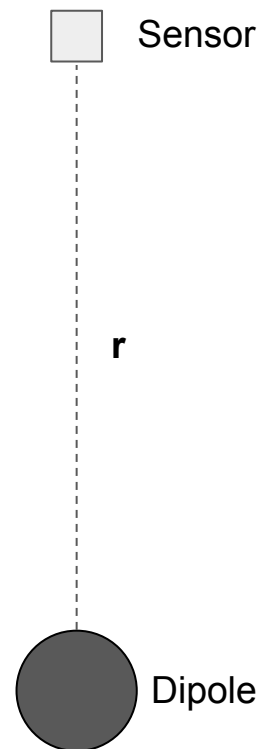
Magnetic Gradient Scale Length Has Been Used in Dipole Localization

For a dipole:

$$\mathbf{r} = -3(\nabla\mathbf{B})^{-1}\mathbf{B}$$

- Useful in RFID localization
 - single dipole-like field
- Cannot use for coil localization
 - multi-coil arrangement and not dipole-like

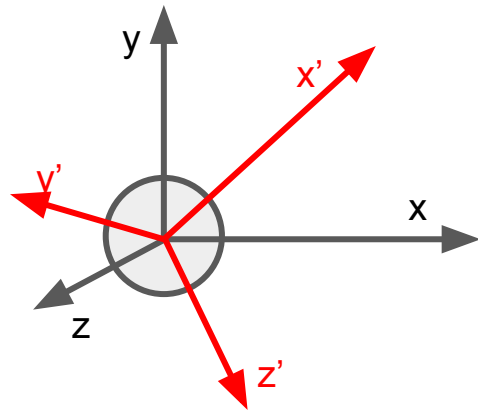
Conclusion: On the right track,
but need a different equation!



Assumptions of $\|\nabla \mathbf{B}\|$ to Formulate $L_{\nabla \mathbf{B}}$

Many matrix norms exist. We should choose one so:

1. Norm uses all 9 components of gradient matrix
2. Norm is invariant to **rotation**



$$\|\mathbf{A}\|_F = \sqrt{\sum_{i=1}^m \sum_{j=1}^n |a_{ij}|^2}$$

$$\nabla \mathbf{B} = \begin{bmatrix} \frac{\partial B_x}{\partial x} & \frac{\partial B_x}{\partial y} & \frac{\partial B_x}{\partial z} \\ \frac{\partial B_y}{\partial x} & \frac{\partial B_y}{\partial y} & \frac{\partial B_y}{\partial z} \\ \frac{\partial B_z}{\partial x} & \frac{\partial B_z}{\partial y} & \frac{\partial B_z}{\partial z} \end{bmatrix}$$

$$L_{\nabla \mathbf{B}} = \frac{\sqrt{2}B}{\|\nabla \mathbf{B}\|_F}$$

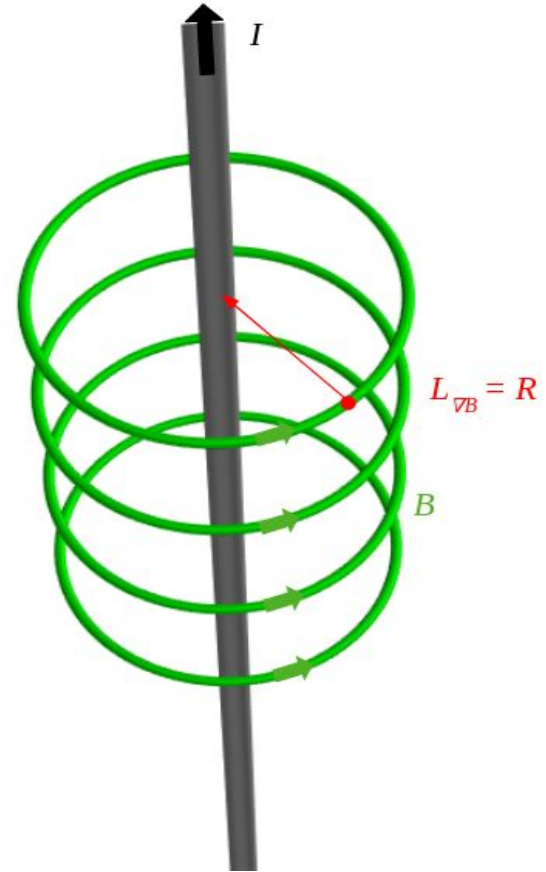
$L_{\nabla\mathbf{B}}$ Behavior In Model Geometry

$$\mathbf{B}(\rho) = \frac{\mu_0 I}{2\pi\rho} \hat{\phi}$$

$$\nabla\mathbf{B} = \frac{\mu_0 I}{2\pi} \left(-\frac{1}{\rho^2} \hat{\phi} \otimes \hat{\rho} - \frac{1}{\rho^2} \hat{\rho} \otimes \hat{\phi} \right)$$

$$\|\nabla\mathbf{B}\|_F = \frac{\sqrt{2}\mu_0 I}{2\pi\rho^2}$$

$$L_{\nabla\mathbf{B}} = \sqrt{2} \frac{B}{\|\nabla\mathbf{B}\|_F} = \rho$$



Outline

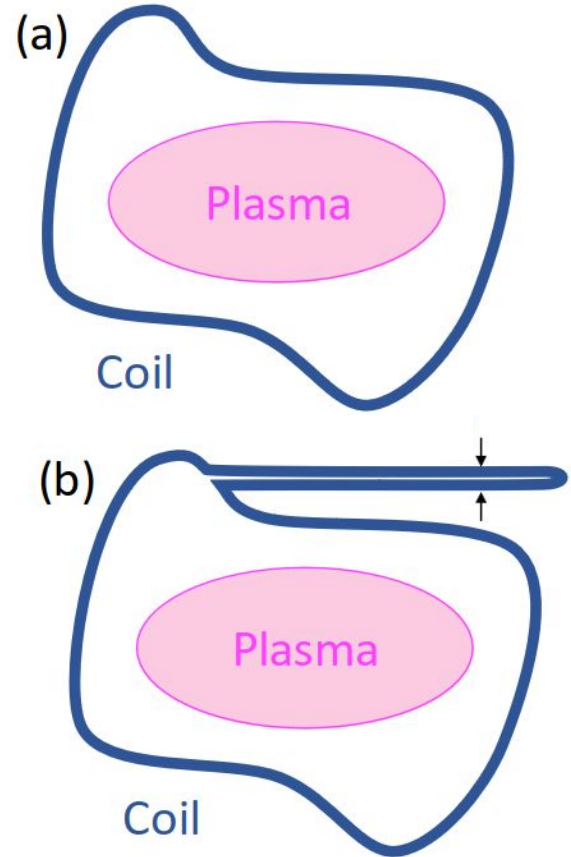
1. Intuition for Magnetic Gradient Scale Length
- 2. Methods of Coil Optimization in REGCOIL**
3. Comparison Between $L_{\nabla B}$ and L_{REGCOIL} and Discussion

REGCOIL is a Useful Optimizer to Systematically Compare the Coils of Many Configurations

For any $\delta > 0$, there exists an infinite number of current arrangements in a finite region around the plasma that match \mathbf{B} on the LCFS to an error $\varepsilon < \delta$.

REGCOIL's objective function **preserves convexity**, so any local minimum is a global minimum.

Fewer tuning parameters than other codes.



Overview of REGCOIL

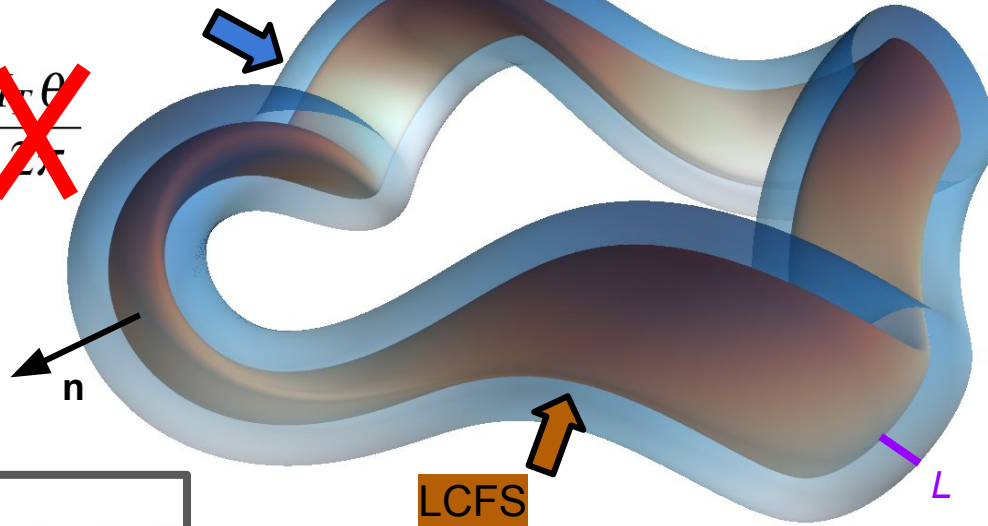
$$\Phi'(\theta', \phi') = \sum_i \Phi_j \sin(m_j \theta' - n_j \phi') + \frac{G\phi'}{2\pi} + \frac{L\theta'}{2\pi}$$

$$\mathbf{K}' = \mathbf{n}' \times \nabla \Phi'$$

$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi} \int d^2 a' \frac{\mathbf{K}' \times (\mathbf{r} - \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3}$$

$$f = \int d^2 a (\mathbf{B}(\theta, \zeta) \cdot \mathbf{n})^2 + \lambda \int d^2 a' \|\mathbf{K}(\theta', \zeta')\|^2$$

Winding Surface



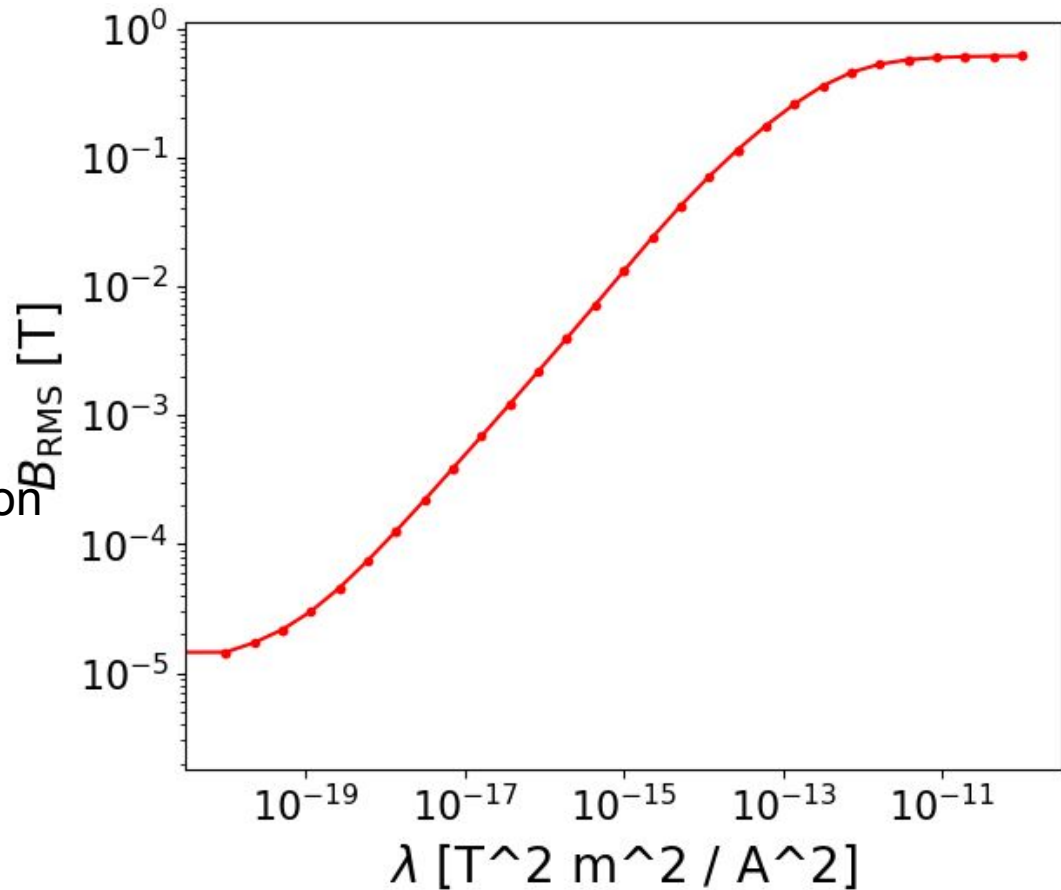
2 free parameters: L and λ . A unique solution requires 2 constraints:

1. $B_{\text{RMS}} = B_{\text{RMS}}^*$
2. $\|K\|_{\infty} = \|K\|_{\infty}^*$

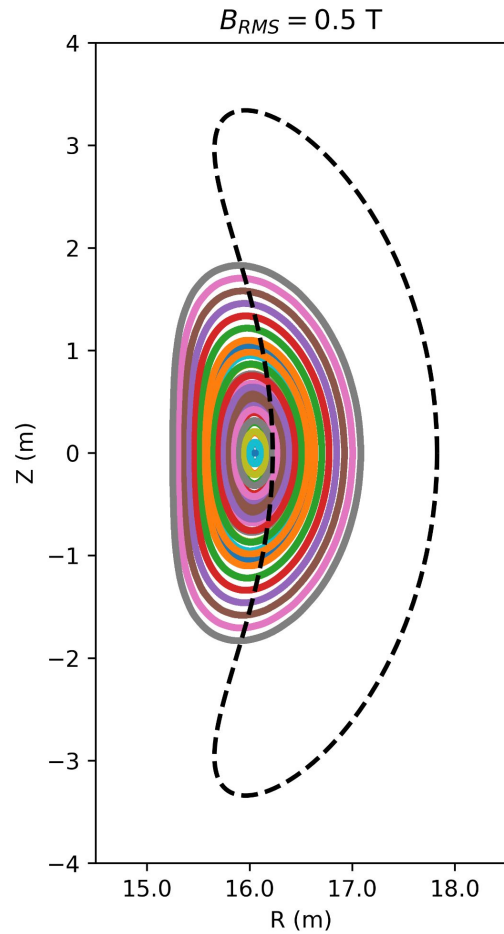
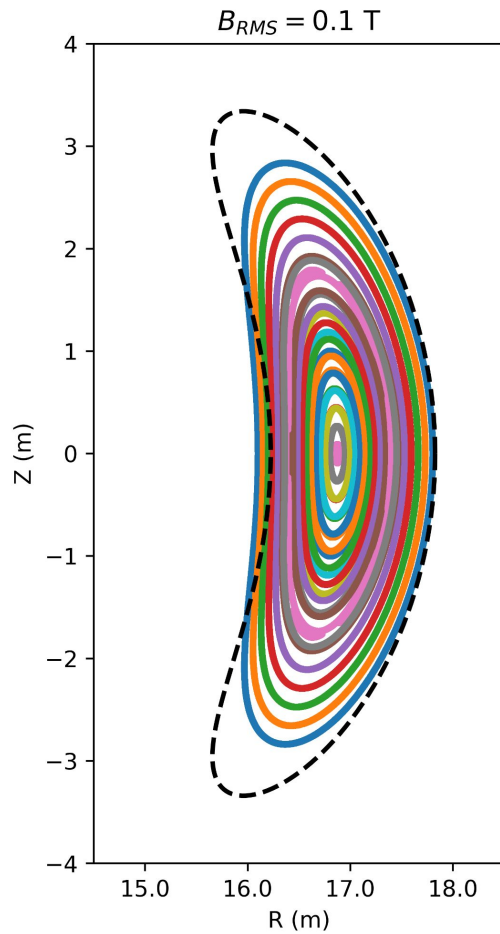
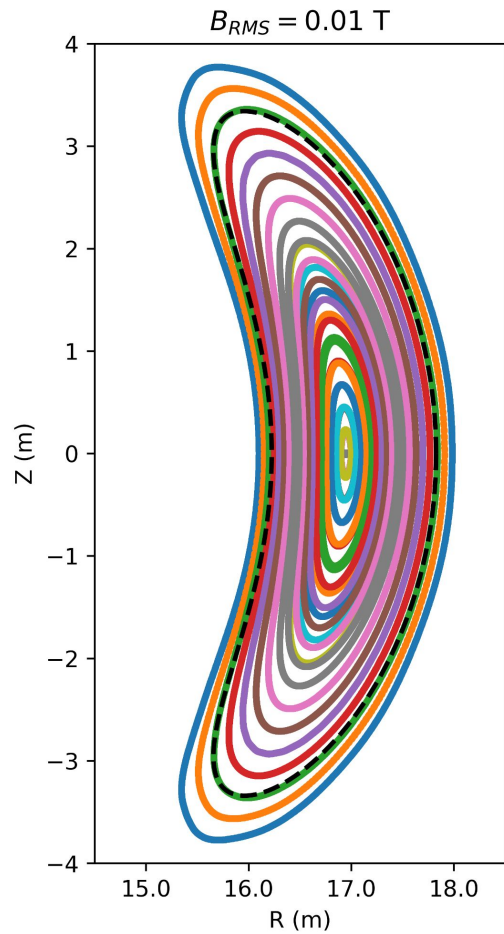
B_{RMS} is a Measure of Accuracy in the Last Closed Magnetic Flux Surface

$$B_{\text{RMS}} = \left(\frac{\int d^2a (\mathbf{B} \cdot \mathbf{n})^2}{A_{\text{Plasma}}} \right)^{1/2}$$

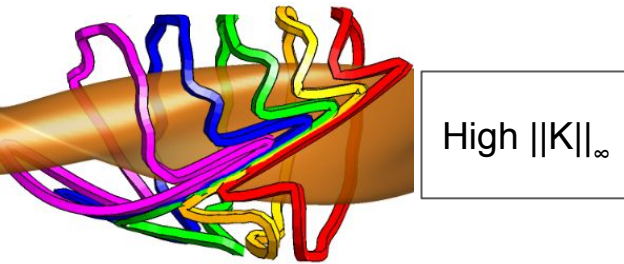
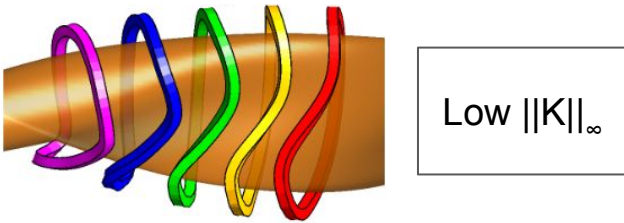
B_{RMS}^* uniquely defines regularization parameter



Smaller B_{RMS} = Better Flux Surfaces

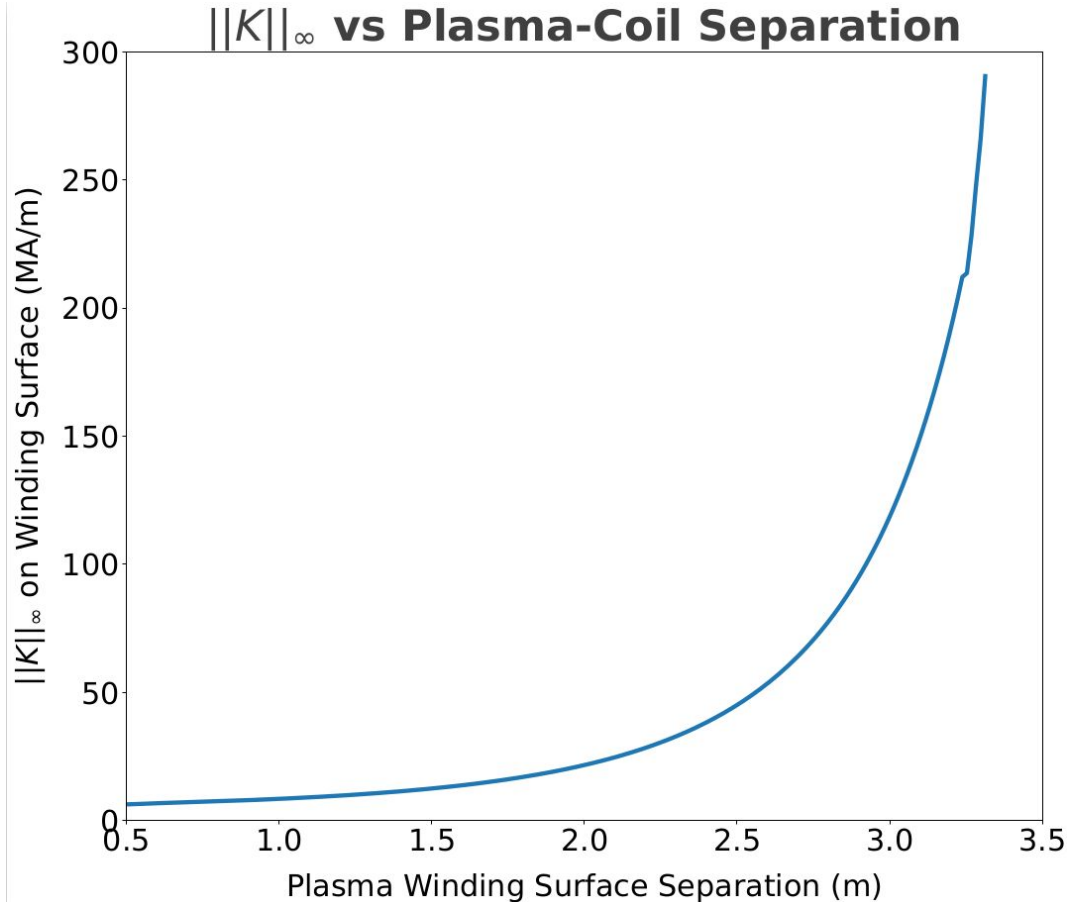


REGCOIL Must be Constrained by $\|K\|_\infty$ to make Buildable Coils



$\|K\|_\infty$ or K_{\max} is the **highest** current density on the winding surface

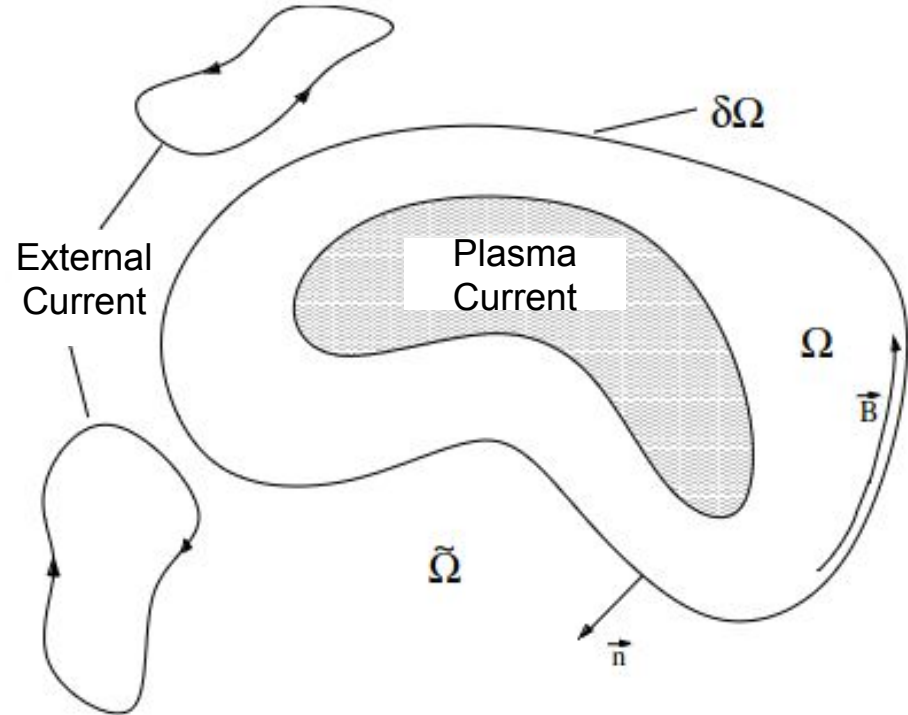
Uniquely defines plasma-coil separation



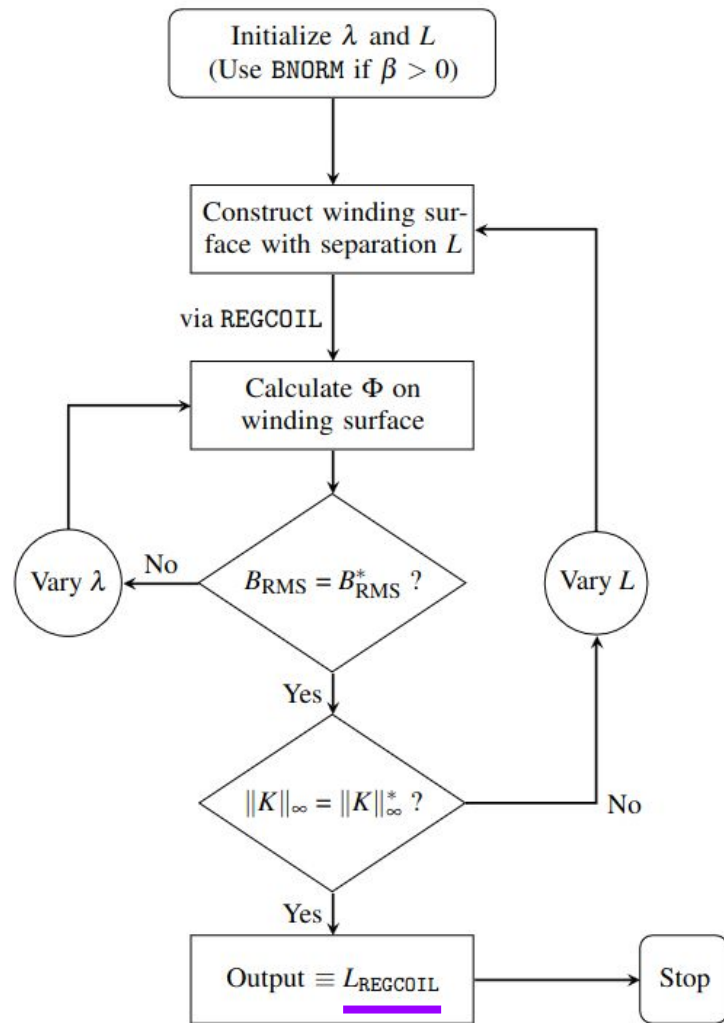
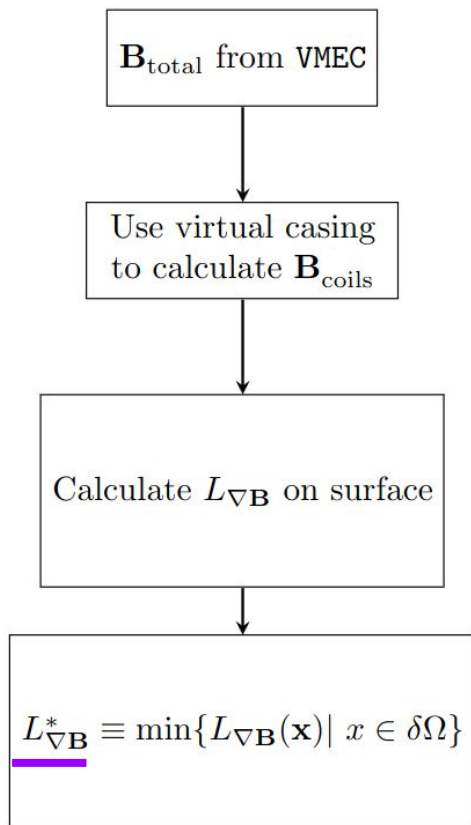
Virtual Casing Decomposes $\mathbf{B}_{\text{coils}}$ From $\mathbf{B}_{\text{total}}$

$$\mathbf{B}(\mathbf{x}) = \mathbf{B}_{\text{plasma}}(\mathbf{x}) + \mathbf{B}_{\text{coils}}(\mathbf{x})$$

$$\mathbf{B}_{\text{coils}}(\mathbf{x}) = -\frac{1}{4\pi} \oint_{\delta\Omega} d^2a \frac{(\mathbf{n} \times \mathbf{B}(\mathbf{p})) \times (\mathbf{x} - \mathbf{p})}{|\mathbf{x} - \mathbf{p}|^3}$$



Summary of Methods



Outline

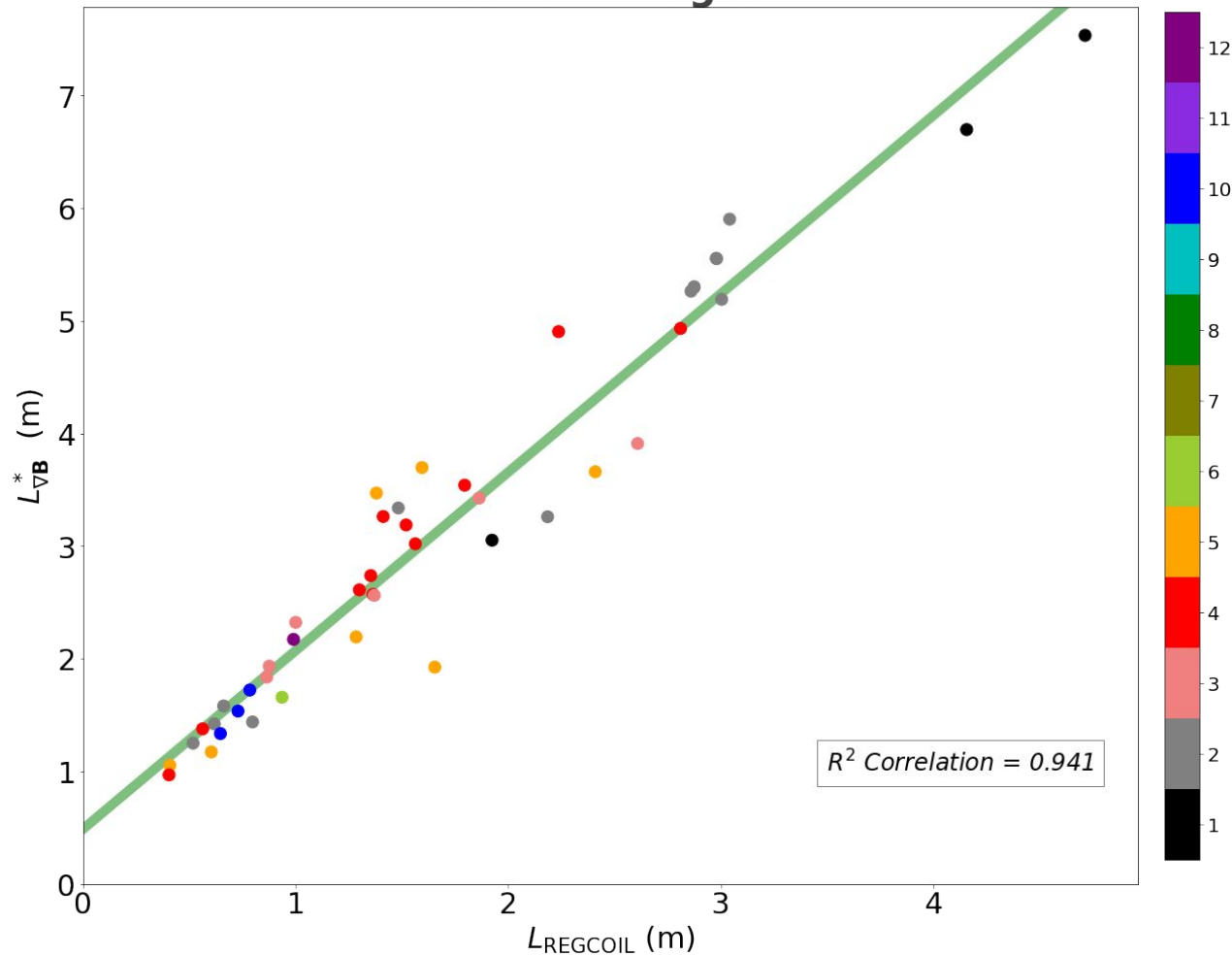
1. Intuition for Magnetic Gradient Scale Length
2. Methods of Coil Optimization in REGCOIL
- 3. Comparison Between $L_{\nabla B}$ and L_{REGCOIL} and Discussion**

$L_{\nabla B}^*$ Accurately Predicts Coil-Plasma Separation Found in Regcoil

N_{fp}

Parameters

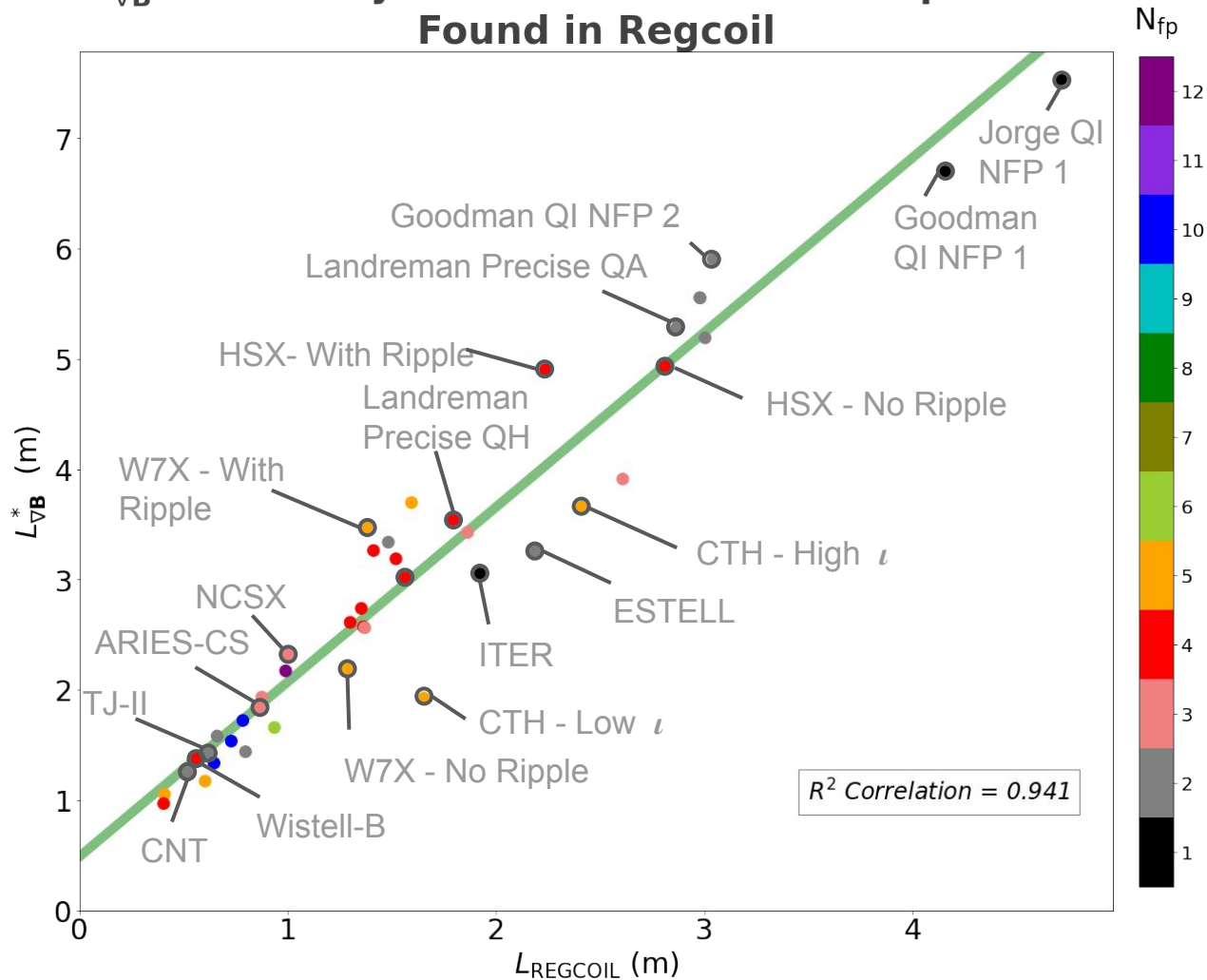
- $B_{RMS}^* = 0.01$ T
- $\|K\|_{\infty}^* = 11.3$ MA/m
- $a = 1.704$ m
- $B_{Vol} = 5.865$ T
- $mpol$ & $ntor = 20$
- n_{θ} & $n_{\zeta} = 96$



$L_{\nabla B}^*$ Accurately Predicts Coil-Plasma Separation Found in Regcoil

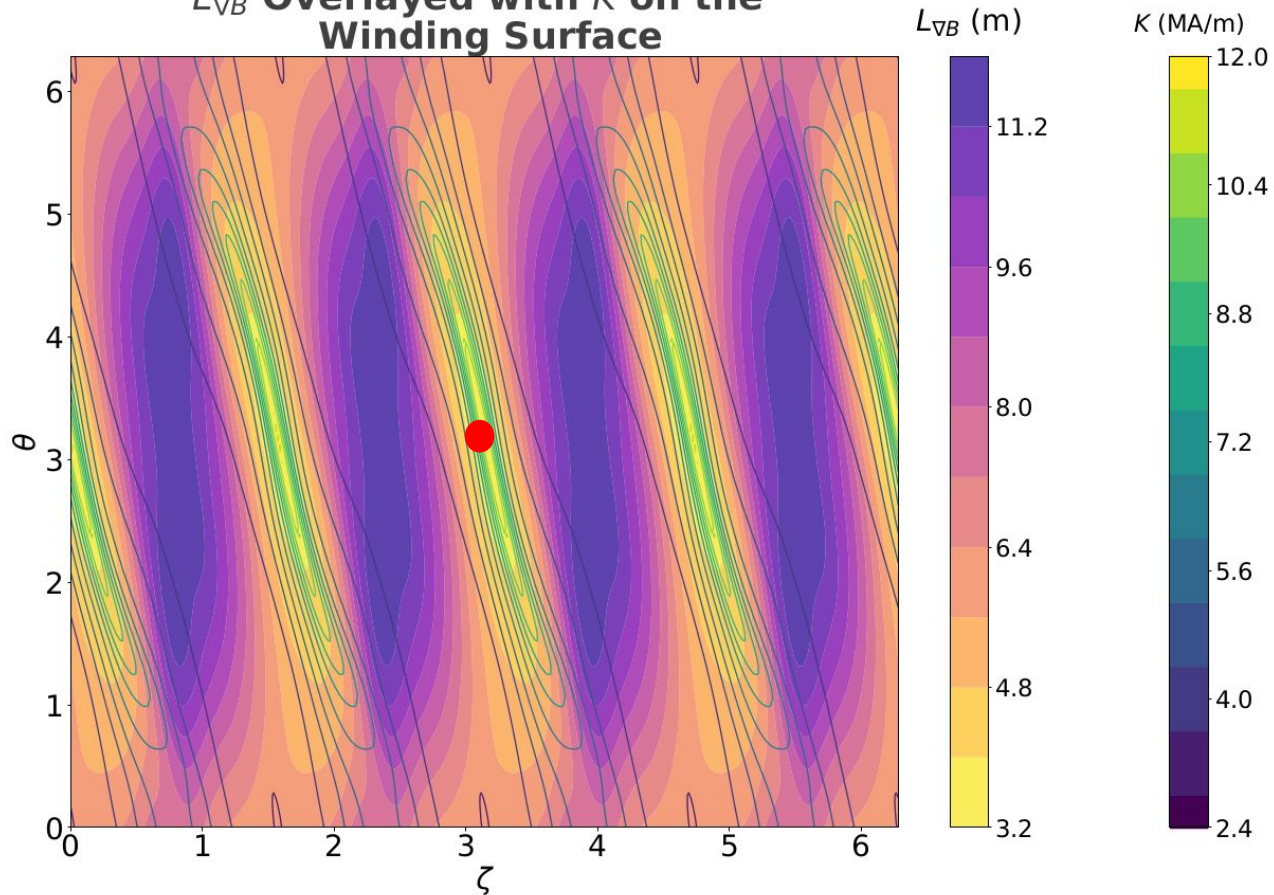
Parameters

- $B_{\text{RMS}}^* = 0.01 \text{ T}$
- $\|K\|_{\infty}^* = 11.3 \text{ MA/m}$
- $a = 1.704 \text{ m}$
- $B_{\text{Vol}} = 5.865 \text{ T}$
- $\text{mpol} \ \& \ \text{ntor} = 20$
- $n_{\theta} \ \& \ n_{\zeta} = 96$



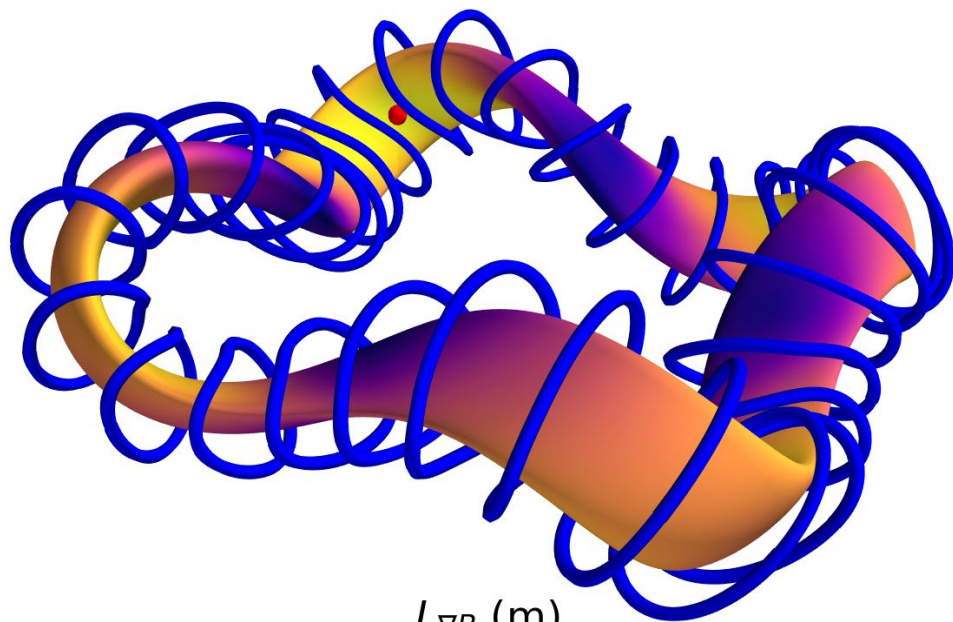
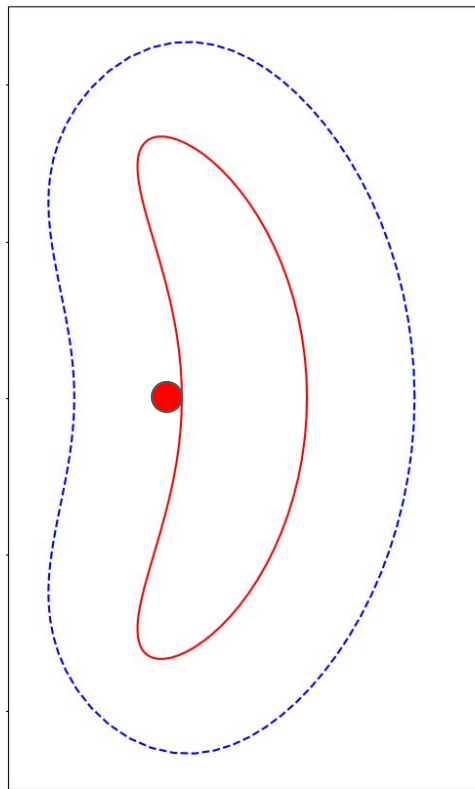
There is Good Spatial Correlation between $\|K\|_\infty$ and $L_{\nabla B}^*$

$L_{\nabla B}$ Overlaid with K on the Winding Surface



The **smallest** $L_{\nabla B}$ and the **largest** K are located at the **same coordinates**

$L_{\nabla B}$ is **Shortest** on the Inside of the Bean Cross-Section



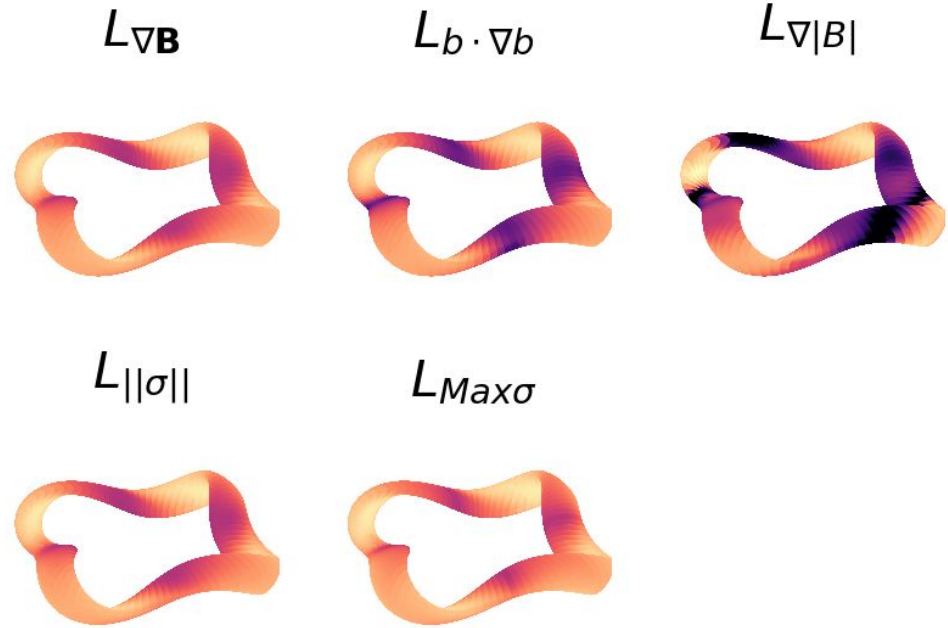
Alternative Magnetic Gradient Scale Lengths are Similar to $L_{\nabla B}$

$$L_{\mathbf{b} \cdot \nabla \mathbf{b}} = \frac{1}{\|\mathbf{b} \cdot \nabla \mathbf{b}\|},$$

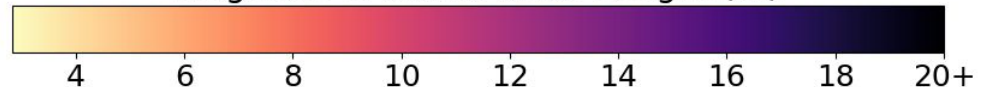
$$L_{\nabla |B|} = \frac{B}{\|\nabla B\|},$$

$$L_{\|\sigma\|} = \frac{\sqrt{2}B}{(\sum_i \sigma_i^2)^{1/2}},$$

$$L_{\text{Max}\sigma} = \frac{B}{\text{Max}[\sigma_i]},$$

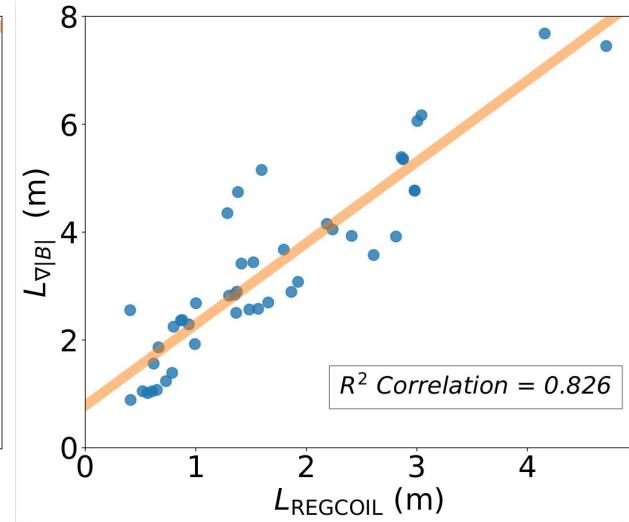
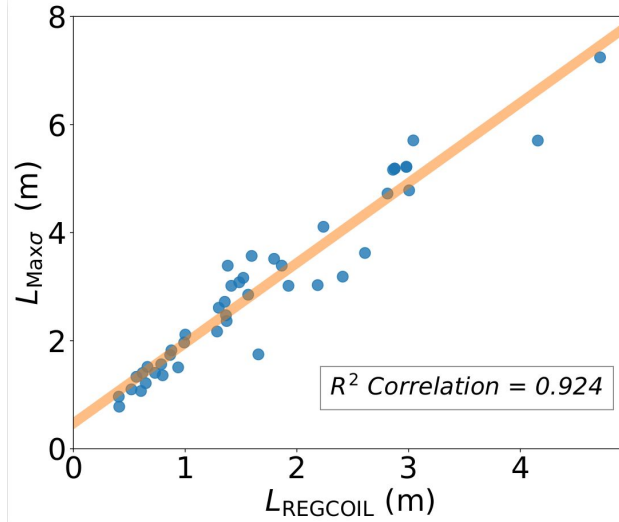
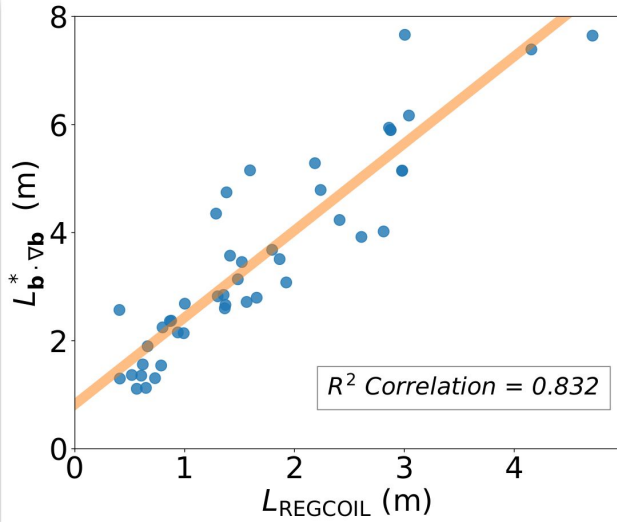


Magnetic Gradient Scale Length (m)



Where σ represents the singular values of $\nabla \mathbf{B}$

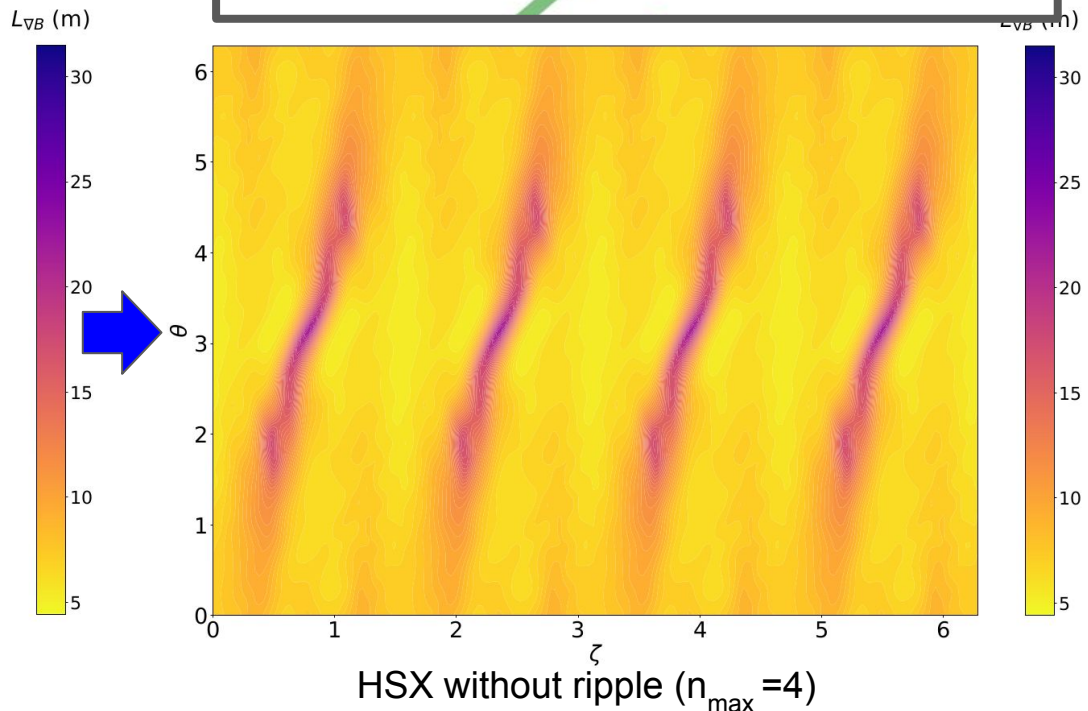
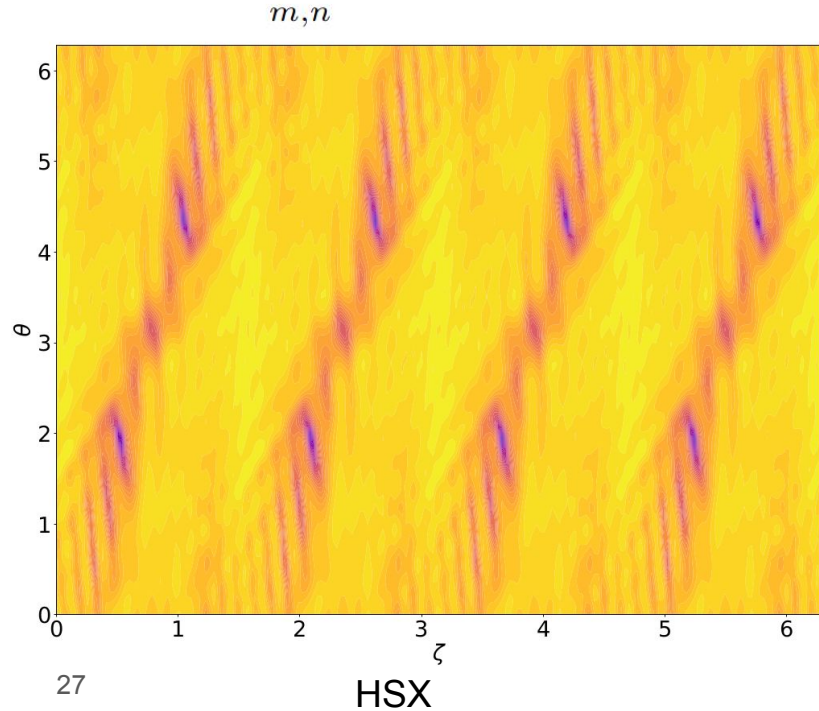
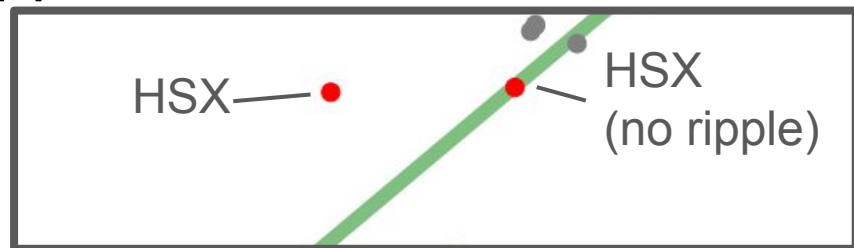
Alternative Scale Lengths



Configurations with High Coil Ripple are Outliers

$$R(\theta, \zeta) = \sum_{m,n} R_{m,n} \cos(m\theta - n\zeta)$$

$$Z(\theta, \zeta) = \sum_{m,n} Z_{m,n} \sin(m\theta - n\zeta)$$



Summary

We established a fundamental connection between $L_{\nabla B}$ and the plasma-coil separation. We calculated the distance of $L_{\nabla B}$ of over 40 configurations, and found a strong correlation between $L_{\nabla B}$ and the plasma-coil separation of stage II optimized configurations with magnetic field accuracy and coil complexity constrained.

$L_{\nabla B}$ is shortest on the inside curve of the bean-shaped cross-section of the plasma, which appears to explain why some stellarators are hard to make with distant coils.

Open Questions/ Ongoing Research

1. Can we get better configurations by optimizing for $L_{\nabla\mathbf{B}}$?
 - a. Currently Ongoing in DESC
2. Is there a better Magnetic Gradient Scale Length than $L_{\nabla\mathbf{B}}$?

$$L_{\nabla\nabla\mathbf{B}} = \frac{4\|\nabla\mathbf{B}\|_F}{\sqrt{2}\|\nabla\nabla\mathbf{B}\|}, \text{ where } \|\nabla\nabla\mathbf{B}\| = \sqrt{\sum_{i,j,k} (\nabla\nabla\mathbf{B})_{ijk}^2}.$$

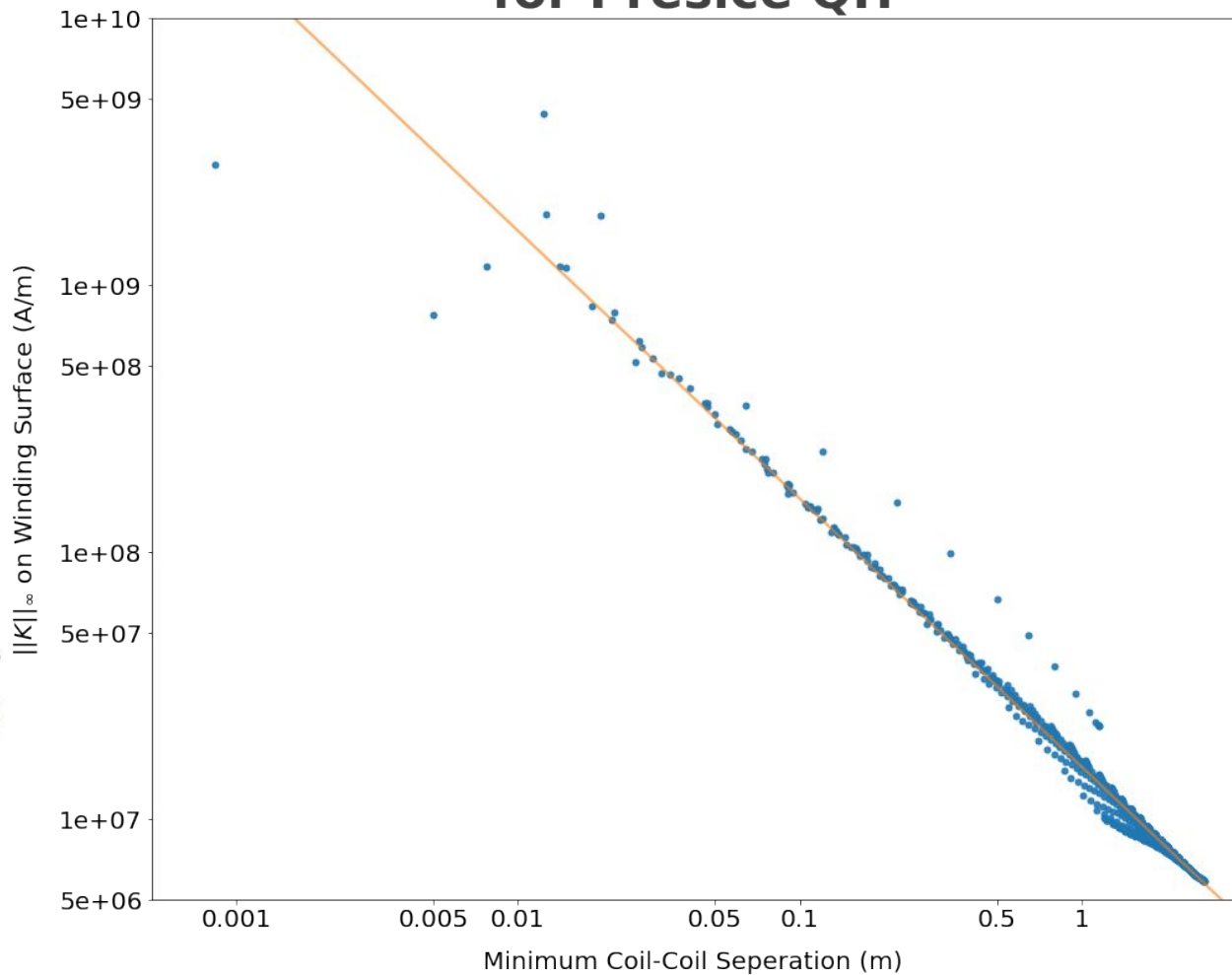
3. How well does $L_{\nabla\mathbf{B}}$ work as a plasma coil separation metric for filamentary coils?
 - a. Can we relax the assumptions that we made in REGCOIL?

$\|K\|_\infty$ is inverse to
Minimum Coil-Coil
Distance

$$\|K\| = \lim_{N \rightarrow \infty} \frac{I_{pol}}{Nd_{cc}}$$

$$K_{\max} \equiv \|K\|_\infty \approx \frac{I_{pol}}{Nd_{min}}$$

Coil Separation vs Current Density for Presice QH



Some VMEC Solutions are Inaccurate in Cartesian Coordinates

Outlier configurations do not pass at least one of the following tests. Most likely caused by computer precision error when converting from VMEC to Cartesian coordinates

$$\frac{\nabla\mathbf{B} - (\nabla\mathbf{B})^T}{2\|\nabla\mathbf{B}\|_F} < 0.38$$

This implies that better accuracy can be achieved by either

a) more accurate VMEC solutions

b) Finding the Frobenius norm without converting to Cartesian coordinates

Full Table of Plasma Configurations (1/2)

Description	N_{fp}	$\beta(\%)$	$L_{REGCDIL}$ (m)	L_{VB}^* (m)
nfp=4 quasi-helical (QH) configuration by Ku & Boozer ^[43]	4	4.00	0.4060	0.9691
Unpublished QH configuration from Michael Drevlak	5	3.92	0.4099	1.0538
Columbia Non-Neutral Torus (CNT) ^[44]	2	0	0.5189	1.2507
Tokamak de la Junta II (TJ-II) ^[45]	4	0	0.5638	1.3777
Wistell-B, Bader et al. ^[46]	5	0	0.6047	1.1726
Quasi-axisymmetric (QA) configuration designed by Paul Garabedian ^[47]	2	3.02	0.6188	1.4214
Large Helical Device (LHD), major radius 3.60m ^[48]	10	0	0.6475	1.3358
Quasi-Poloidal Stellarator (QPS) ^[49]	2	2.01	0.6625	1.5812
LHD, major radius 3.53m ^[50]	10	0	0.7302	1.5354
LHD, major radius 3.75m ^[51]	10	0	0.7858	1.7226
Henneberg et al. QA ^[50]	2	3.50	0.7987	1.4390
Advanced Research Innovation and Evaluation Study-Compact Stellarator (ARIES-CS) ^[1]	3	4.06	0.8655	1.8375
National Compact Stellarator Experiment (NCSX) stage-1 optimization result (known as LI383) ^[52]	3	4.26	0.8771	1.9343
The first quasisymmetric configuration found ^[53]	6	4.09	0.9374	1.6582
Advanced Toroidal Facility (ATF) ^[54]	12	4.48	0.9913	2.1721
NCSX free-boundary (c09r00) ^[55]	3	4.08	1.0015	2.3233
Wendelstein 7-X (W7-X), without coil ripple ^[56]	5	4.48	1.2858	2.1941
Landreman, Buller, & Drevlak, QH, 5% beta ^[57]	4	5.58	1.3009	2.6120
Landreman, Buller, & Drevlak, QH, vacuum ^[58]	4	0	1.3545	2.7385
Boundary constructed by near-axis expansion. Vacuum QH with nfp=4	3	0	1.3650	2.5722
Goodman et al. Quasi-isodynamic (QI) configuration with nfp=3 ^[59]	4	0	1.3712	2.5633

Full Table of Plasma Configurations (2/2)

Quasi-Isodynamic (QI) configuration from CIEMAT ^[28]	4	0	1.4130	3.2634
Chinese First Quasiaxisymmetric Stellarator (CQFS) ^[59]	2	0	1.4839	3.3392
Landreman & Paul, QH with magnetic well ^[27]	4	0	1.5206	3.1882
Wistell-A. Bader et al. ^[60]	4	0	1.5641	3.0210
W7-X "high narrow mirror" configuration ^[61]	5	4.00	1.5952	3.6979
Compact Toroidal Hybrid (CTH) Stellarator, vacuum, with low rotational transform ^[41]	5	0	1.6556	1.9259
Landreman & Paul, precise QH ^[27]	4	0	1.7960	3.5418
Unpublished nfp=3 QH	3	0	1.8644	3.4277
Up-down-symmetric ITER-like configuration ^[39]	1	2.28	1.9248	3.0531
Evolutionary Stellarator of Lorraine (ESTELL) ^[62]	2	0	2.1860	3.2610
Helically Symmetric Experiment (HSX), standard configuration, vacuum, with coil ripple ^[40]	4	0	2.2377	4.9052
Compact Toroidal Hybrid (CTH) stellarator, vacuum, with high rotational transform ^[41]	5	0	2.4102	3.6607
Boundary constructed by near-axis expansion. Vacuum QH with nfp=3 ^[24]	3	0	2.6091	3.9118
HSX, standard configuration, vacuum, without coil ripple	4	0	2.8111	4.9336
Vacuum QA configuration with 16 coils from Giuliani et al. Coil length 24m. ^[63]	2	0	2.8602	5.2643
Landreman & Paul, precise QA ^[27]	2	0	2.8748	5.2977
Wechsung et al. QA without magnetic well, coil length 24m. ^[64]	2	0	2.8750	5.3037
Wechsung et al. QA with magnetic well, coil length 24m ^[64]	2	0	2.9790	5.5563
Landreman & Paul QA with magnetic well. ^[27]	2	0	2.9806	5.5532
Goodman et al. Quasi-isodynamic configuration with nfp=2 ^[56]	2	0	3.0045	5.1919
Landreman, Buller & Drevlak, QA, 2.5% beta ^[55]	2	2.55	3.0419	5.9042
Goodman et al. Quasi-isodynamic configuration with nfp=1 ^[56]	1	0	4.1563	6.6993
Jorge et al. Quasi-isodynamic configuration with nfp=1 ^[65]	1	0	4.7133	7.5360

$L_{\nabla\mathbf{B}}$ Behavior In A Circular Wire

$$\rho^2 = x^2 + y^2; r^2 = x^2 + y^2 + z^2; \alpha^2 = 1 + r^2 - 2\rho$$

$$\beta^2 = 1 + r^2 + 2\rho; k^2 = 1 - \frac{\alpha^2}{\beta^2}$$

$$B_x = \frac{x}{2\alpha^2\beta\rho^2} [(1+r^2)E(k^2) - \alpha^2K(k^2)]$$

$$B_y = \frac{y}{2\alpha^2\beta\rho^2} [(1+r^2)E(k^2) - \alpha^2K(k^2)]$$

$$B_z = \frac{1}{2\alpha^2\beta} [(1-r^2)E(k^2) + \alpha^2K(k^2)],$$

$L_{\nabla\mathbf{B}}$ for a Magnetic Field of a Circular Wire

